

DUALITY



SUPERSYMMETRY

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- DUALITY SYMMETRIES OF THE LOW-ENERGY STRING EFFECTIVE ACTION:

1. - $SL(2, R)$ (S) - DUALITY
2. - TARGET-SPACE (R) - DUALITY

- SUPERSYMMETRIC BOSONIC CONFIGURATIONS

- KILLING SPINORS
- BOGOMOLNYI BOUNDS - EXTREMALITY

- THE EFFECT OF S- AND R-DUALITY ON THE SUPERSYMMETRY PROPERTIES

THE LOW-ENERGY STRING EFFECTIVE ACTION

$$S = \frac{1}{2} \int d^D x \sqrt{g} e^{-2\phi} \left\{ -R + 4(\partial\phi)^2 - \frac{3}{4} H^2 + O(\alpha') \right\}$$

$$H_{\mu\nu\rho} = \partial_{[\mu} B_{\nu\rho]} \quad B_{\nu\rho} \rightarrow \text{axion 2-form}$$

STARTING POINT $D=10$ ($N=1, d=10$ SUGRA)



$D=4$ (+ 6 Abelian vector fields)
 $N=4, d=4$ SUGRA

$$H_{\mu\nu\rho} = \partial_{[\mu} B_{\nu\rho]} + \text{Chern-Simons}$$

(Only) in $D=4$ one can perform the change of variables

$$H_{\mu\nu\rho} \sim \epsilon_{\mu\nu\rho\sigma} \partial_\sigma a \leftarrow \text{pseudoscalar axion}$$

$$\lambda = a - i e^{-2\phi}; \quad \tilde{F} = \lambda F^+ - \bar{\lambda} F^-$$

$$g_{\mu\nu}^{\text{can}} = e^{-2\phi} g_{\mu\nu}^{\text{string}}$$

$$S = \frac{1}{2} \int d^4 x \sqrt{g} \left\{ -R + 2 \frac{\partial\lambda\partial\bar{\lambda}}{(\text{Im}\lambda)^2} - F\tilde{F} \right\} \quad \begin{matrix} N=4 \\ d=4 \\ \text{SUGRA} \end{matrix}$$

Eq of motion $\nabla_\mu * \tilde{F}^{\mu\nu} = 0;$
 Bianchi identity $\nabla_\mu * F^{\mu\nu} = 0;$ duality 3

$$\begin{cases} \tilde{F}' = \alpha \tilde{F} + \beta F; \\ F' = \gamma \tilde{F} + \delta F; \end{cases} \quad \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \in SL(2, \mathbb{R})$$

$$\Rightarrow \lambda' = \frac{\alpha \lambda + \beta}{\gamma \lambda + \delta};$$

$SL(2, \mathbb{R})$ DOES NOT ACT ON THE CANONICAL METRIC

BUT IT DOES ON THE STRING METRIC

IT CAN BE USED TO GENERATE NEW SOLUTIONS:

MANIFESTLY $SL(2, \mathbb{R})$ -INVARIANT FAMILY OF DILATON-

AXION TAUB-NUT METRICS

(KASTOR, TORMA, GALT'SOV, GARCÍA MYERS, JOHNSON)

$$ds^2 = f(dt + 2\ell \omega \theta d\varphi)^2 - f^{-1} dr^2 - R^2 d\Omega^2;$$

$$f = \frac{(\tilde{r} - \tilde{r}_+)(\tilde{r} - \tilde{r}_-)}{R^2}; \quad R^2 = r^2 + \ell^2 - |\gamma|^2;$$

$$\tilde{r}_\pm = m \pm \tilde{r}_0; \quad \tilde{r}_0^2 = m^2 + \ell^2 + |\gamma|^2 - 4|\tilde{r}|^2;$$

$$\lambda = \frac{\tilde{r}_0(\tilde{r} + i\ell) + \tilde{r}_0 \gamma}{(\tilde{r} + i\ell) + \gamma}; \quad \gamma = \frac{-2\tilde{r}^2}{M} \dots$$

(4)

- $SL(2, R)$ IS QUANTUM-MECHANICALLY
BROKEN TO $SL(2, Z)$ - CHARGE QUANTIZATION
- INSTANTONS

- $SL(2, R)$ INTERCHANGES THE STRONG-
AND WEAK-COUPLING REGIMES OF
STRING THEORY $\phi \rightarrow -\phi$

- $SL(2, Z)$ COULD BE AN EXACT SYMMETRY
OF STRING THEORY (FONT, IBÁÑEZ, LÜST, QUEVEDO
REY
SEN)

- $SL(2, Z)$ - INVARIANT SPECTRUM (SEN, SCHWARZ, KALLOSH
...)

- STATES = BLACK HOLES? (DUFF)

TARGET-SPACE DUALITY

(5)

FOR ANY D , IF A SOLUTION OF THE EQUATIONS OF MOTION $\{g_{\mu\nu}, B_{\mu\nu}, \phi\}$ DOES NOT DEPEND ON THE COORDINATE x , THERE EXISTS A DUAL CONFIGURATION $\{\tilde{g}_{\mu\nu}, \tilde{B}_{\mu\nu}, \tilde{\phi}\}$

$$\tilde{g}_{xx} = 1/g_{xx}; \quad \tilde{g}_{x\mu} = B_{x\mu}/g_{xx};$$

$$\tilde{g}_{\mu\nu} = g_{\mu\nu} - (g_{x\mu}g_{x\nu} - B_{x\mu}B_{x\nu})/g_{xx};$$

$$\tilde{B}_{x\mu} = g_{x\mu}/g_{xx}; \quad \tilde{B}_{\mu\nu} = B_{\mu\nu} + (g_{x\mu}B_{\nu x} - g_{x\nu}B_{\mu x})/g_{xx}$$

$$\tilde{\phi} = \phi - \frac{1}{2} \ln |g_{xx}| \quad (\text{BUSCHER})$$

WHICH IS ALSO A SOLUTION OF THE EQS. OF MOTION.

FROM THE σ -MODEL POINT OF VIEW BOTH CONFIGURATIONS CORRESPOND TO THE SAME THEORY

EXAMPLE: $\mathbb{R} \longrightarrow 1/\mathbb{R}$

IT IS THE MOST CHARACTERISTIC SYMMETRY OF STRING THEORY

SUPERSYMMETRIC BOSONIC CONFIGURATIONS

- THEORY WITH LOCAL SUSY $\{\phi^B, \phi^F\}$
- PARTICULAR SOLUTIONS BREAK SUSY.
- CONSIDER CONFIGURATIONS $\{\phi^B, \phi^F=0\}$.
CAN THEY BE INVARIANT UNDER SUSY?

$$\delta_\epsilon \phi^B \propto \phi^F = 0 \quad \checkmark$$

$$\delta_\epsilon \phi^F \propto \phi^B \neq 0 \quad \text{IN GENERAL}$$

$$\delta_\epsilon \phi^F \equiv D_\mu \epsilon = 0 \quad \text{KILLING SPINOR EQUATION}$$

- THE EXISTENCE OF KILLING SPINORS IS ASSOCIATED WITH THE SATURATION OF BOGOMOLNY BOUNDS AND (SOMETIMES) WITH EXTREMALITY.
- SUPERSYMMETRIC CONFIGURATIONS ARE FREE OF QUANTUM CORRECTIONS.

IS SUSY PRESERVED BY S-OR-R-DUALITY?

SUSY vs. S-DUALITY

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TH: IF $\{g_{\mu\nu}, \lambda, F_{\mu\nu}\}$ ADMITS A KILLING SPINOR ϵ ,
ANY S-DUAL CONFIGURATION $\{g_{\mu\nu}, \lambda', F'_{\mu\nu}\}$
ALSO ADMITS A KILLING SPINOR

$$\epsilon' \sim e^{i \int \text{Im}(\lambda' \delta)} \epsilon \quad \begin{pmatrix} \alpha & \beta \\ 0 & \delta \end{pmatrix} \in \text{SL}(2, \mathbb{R})$$

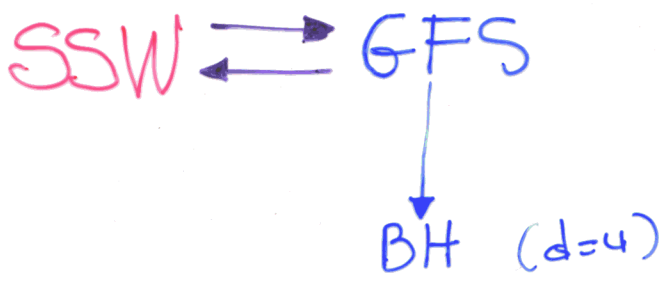
→ THE BOGOMOLNYI BOUND $M^2 + |\rho|^2 - 4|\tau|^2 > 0$
IS INVARIANT UNDER S-DUALITY (T.O. SEN)

→ THIS PROPERTY IS CONSISTENT WITH
THE EXISTENCE OF A SUPERSYMMETRIC
 $\text{SL}(2, \mathbb{Z})$ -INVARIANT SPECTRUM OF THE
HETEROTIC STRING → **EXTREME
BLACK
HOLES?**

SUSY vs. R-DUALITY

TH: IF $\{g_{\mu\nu}, B_{\mu\nu}, \phi\}$ IS A BOSONIC CONFIGURATION WHICH DOES NOT DEPEND ON THE COORDINATE x AND ADMITS AN $N=1, d=10$ KILLING SPINOR ϵ WHICH DOES NOT DEPEND ON x , EITHER, THEN

- 1) IF x IS SPACE-LIKE, ϵ IS ALSO KILLING FOR THE DUAL CONFIGURATION
- 2) IF x IS TIME-LIKE, $\tilde{\epsilon} = \gamma_x \epsilon$ IS ALSO KILLING FOR THE DUAL CONFIGURATION.



BERGSHOFF
ENTROP
KALLOSU

DUALITY @ SUSY IS GOVERNED BY AN EFFECTIVE $d=9$ -DIMENSIONAL THEORY

FUTURE DIRECTIONS

- INCLUSION OF VECTOR FIELDS (α')
 - EXTENSION TO NON-ABELIAN DUALITY
 - $N=4, d=4$ CASE
- COSMIC CENSORSHIP VS. SUSY
- DUALITY IN SUPERSPACE

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